

Reduced-Order Models for Thermal Radiative Transfer Based on POD-Galerkin Method and Low-Order Quasidiffusion Equations

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- We consider problems of high-energy density physics, where radiative transfer is the main mechanism of energy redistribution
 - Described by complex systems of multiphysical differential equations
- Numerical simulations of these complex multiphysical problems are faced with several challenges
 - Strong nonlinearity
 - Multi-scale characterization
 - **High-dimensionality**
- The Boltzmann transport equation (BTE) drives the dimensionality
 - The solution is 7-dimensional
 - Independent variables include: time (t), spatial position (\mathbf{r}), direction of motion ($\mathbf{\Omega}$), frequency (ν) describing photons
 - Simple discretization - 100 nodes each axis: 10^{12} degrees of freedom
 - The high dimensionality imposes large computational burden and memory footprint
- Reduced order models for the BTE are commonly employed to reduce the problem of dimensionality

- We develop a reduced-order model (ROM) for the BTE to reduce the dimensionality of multiphysical simulations
 - Reducing both computational cost and memory requirements
- Basic idea of the model:
 - Formulate a proper orthogonal decomposition - Galerkin (POD-G) projection for the BTE using known solution data
 - The photon intensities are expanded about the POD basis
 - The projected form of the BTE solves for the coefficients of this expansion for photon intensities
 - Equip the low-dimensional projected BTE with a system of low-order moment equations of the BTE
- The POD-G BTE serves to calculate closures for the moment equations, which are coupled to the specific multiphysics equations of interest

- The first prototype of this method is formulated on the multigroup thermal radiative transfer (TRT) problem in 1D slab geometry
- The high-order Boltzmann transport equation

$$\frac{1}{c} \frac{\partial I_g}{\partial t}(x, \mu, t) + \mu \frac{\partial I_g}{\partial X}(x, \mu, t) + \kappa_g(T) I_g(x, \mu, t) = 2\pi \kappa_g(T) B_g(T) \quad (1)$$

$$x \in [0, X], \quad \mu \in [-1, 1], \quad g = 1, \dots, N_g, \quad t \geq 0,$$

$$I_g|_{\substack{\mu>0 \\ x=0}} = I_g^{in+}, \quad I_g|_{\substack{\mu<0 \\ x=X}} = I_g^{in-}, \quad I_g|_{t=0} = I_g^0, \quad (2)$$

- The material energy balance equation

$$\frac{\partial \varepsilon(T)}{\partial t} = \sum_{g=1}^{N_g} \kappa_g(T) \left(\int_{-1}^1 I_g(x, \mu, t) d\mu - 4\pi B_g(T) \right), \quad T|_{t=0} = T_0. \quad (3)$$

- Supersonic radiation flow problem neglecting material motion, photon scattering, heat conduction and external sources

- The POD-G projection method is formulated in discrete space
- The high-order Boltzmann transport equation

$$\frac{1}{c} \frac{\partial I_g}{\partial t}(x, \mu, t) + \mu \frac{\partial I_g}{\partial x}(x, \mu, t) + \kappa_g(T) I_g(x, \mu, t) = 2\pi \kappa_g(T) B_g(T)$$

- Discretize with: Discrete-Ordinates, Backward-Euler, Simple Corner Balance

$$\frac{1}{c \Delta t^n} (\mathbf{I}^n - \mathbf{I}^{n-1}) + \mathcal{L}_h \mathbf{I}^n + \mathcal{K}_h^n(T) \mathbf{I}^n = \mathbf{Q}^n(T), \quad (4)$$

- Discrete operators $\mathcal{L}_h, \mathcal{K}_h^n(T)$ determined by scheme
- N_x spatial cells, N_μ discrete directions, N_t time steps,
- $D = 2N_x N_\mu N_g$
- Solution vector: $\mathbf{I}^n = ((\mathbf{I}_1^n)^\top \dots (\mathbf{I}_{N_g}^n)^\top)^\top \in \mathbb{R}^D$
- Construct snapshot matrix

$$\mathbf{A} = [\mathbf{I}^1, \dots, \mathbf{I}^{N_t}] \quad (5)$$

- Goal: expand intensities in basis functions $\{\mathbf{u}_\ell\}_{\ell=1}^r$

$$\mathbf{l}_r^u(t^n) = \sum_{\ell=1}^r \lambda_\ell^n \mathbf{u}_\ell \quad (6)$$

- We formulate the POD basis $\{\mathbf{u}_\ell\}_{\ell=1}^r$, $r \ll D$ using snapshots in \mathbf{A}

$$\min_{\mathbf{u}_1, \dots, \mathbf{u}_r} \sum_{n=1}^{N_t} \Delta t^n \left\| \mathbf{l}^n - \sum_{\ell=1}^r \langle \mathbf{l}^n, \mathbf{u}_\ell \rangle_W \mathbf{u}_\ell \right\|_W^2, \quad (7)$$

- Weighted inner product specific to the discretization: $\langle \mathbf{u}, \mathbf{v} \rangle_W = \mathbf{u}^\top \mathbf{W} \mathbf{v}$

- Standard POD uses the identity matrix $\mathbf{W} = \mathbb{I}$ so that $\langle \mathbf{u}, \mathbf{v} \rangle_{\mathbf{W}} = \langle \mathbf{u}, \mathbf{v} \rangle$
- We seek \mathbf{W} to correspond to the discrete integration over space, angle, frequency
- For the considered discretization schemes we have

$$\int_0^{\infty} \int_{-1}^1 \int_0^{L_x} u(x, \mu, \nu) dx d\mu d\nu \Rightarrow \sum_{g=1}^{N_g} \sum_{m=1}^{N_{\mu}} w_m \sum_{i=1}^{N_x} \frac{\Delta x_i}{2} (\mathbf{u}_{g,m,i,L} + \mathbf{u}_{g,m,i,R}) \quad (8)$$

- We find the matrix \mathbf{W} as

$$\mathbf{W} = \bigoplus_{g=1}^{N_g} \bigoplus_{m=1}^{N_{\mu}} w_m \hat{\mathbf{W}}^x, \quad \hat{\mathbf{W}}^x = \bigoplus_{i=1}^{N_x} \begin{pmatrix} \frac{\Delta x_i}{2} & 0 \\ 0 & \frac{\Delta x_i}{2} \end{pmatrix} \quad (9)$$

- Construct snapshot matrix

$$\mathbf{A} = [\mathbf{I}^1, \dots, \mathbf{I}^{N_t}] \quad (10)$$

- Calculate weighted snapshot matrix

$$\hat{\mathbf{A}} = \mathbf{W}^{1/2} \mathbf{A} \mathbf{D}^{1/2}, \quad \mathbf{D} = \text{diag}(\Delta t^1, \dots, \Delta t^{N_t}) \quad (11)$$

- Find singular value decomposition of $\hat{\mathbf{A}}$

$$\hat{\mathbf{A}} = \hat{\mathbf{U}} \hat{\mathbf{S}} \hat{\mathbf{V}}^T \quad (12)$$

$$\hat{\mathbf{U}} = [\hat{\mathbf{u}}_1, \dots, \hat{\mathbf{u}}_d], \quad \hat{\mathbf{S}} = \text{diag}(\sigma_1, \dots, \sigma_d), \quad \hat{\mathbf{V}} = [\hat{\mathbf{v}}_1, \dots, \hat{\mathbf{v}}_d] \quad (13)$$

- The POD basis is then found as $\mathbf{U} = [\mathbf{u}_1, \dots, \mathbf{u}_d]$ with $d = \text{rank}(\hat{\mathbf{A}})$ using

$$\mathbf{U} = \mathbf{W}^{-1/2} \hat{\mathbf{U}} \quad (14)$$

- Full rank system of equations for $\{\lambda_\ell^n\}$ is of size d using all \mathbf{u}_ℓ
- We seek a system of size $r \leq d$ using the first r basis functions in the expansion
- As r decreases, computational efficiency will increase while accuracy decreases
- We seek $r \ll D$ that will give certain level of accuracy
- Truncation criteria (tuning parameter):

$$\xi^2 = \frac{\sum_{\ell=r+1}^d \sigma_\ell^2}{\sum_{\ell=1}^d \sigma_\ell^2} \quad (15)$$

- Set some ξ and find rank r that satisfies the relation above
- This allows the method to easily trade computational requirements with accuracy at will per simulation

$$\frac{1}{c\Delta t^n} (\mathbf{I}^n - \mathbf{I}^{n-1}) + \mathcal{L}_h \mathbf{I}^n + \mathcal{K}_h^n(T) \mathbf{I}^n = \mathbf{Q}^n(T), \quad \mathbf{I}_r^u(t^n) = \sum_{\ell=1}^r \lambda_\ell^n \mathbf{u}_\ell$$

- POD Galerkin-Projected BTE (apply $\langle \mathbf{u}_\ell, \cdot \rangle_W$)

$$\begin{aligned} \frac{1}{c\Delta t^n} (\lambda_\ell^n - \lambda_\ell^{n-1}) + \sum_{\ell'=1}^r \lambda_{\ell'}^n \langle \mathbf{u}_\ell, \mathcal{L}_h \mathbf{u}_{\ell'} \rangle_W \\ + \sum_{\ell'=1}^r \lambda_{\ell'}^n \langle \mathbf{u}_\ell, \mathcal{K}_h^n(T) \mathbf{u}_{\ell'} \rangle_W = \langle \mathbf{u}_\ell, \mathbf{Q}^n(T) \rangle_W \quad (16) \end{aligned}$$

- Used orthogonality of basis: $\langle \mathbf{u}_{\ell'}, \mathbf{u}_\ell \rangle_W = \delta_{\ell, \ell'}$
- Dense system of equations for $\{\lambda_\ell^n\}$

- High-order Boltzmann transport equation

$$\frac{1}{c} \frac{\partial I_g}{\partial t} + \mu \frac{\partial I_g}{\partial x} + \kappa_g(T) I_g = 2\pi \kappa_g(T) B_g(T)$$

- Eddington factor $f_g[l] = \int_{-1}^1 \mu^2 I_g d\mu / \int_{-1}^1 I_g d\mu$

- Multigroup quasidiffusion equations for $E_g = \frac{1}{c} \int_{-1}^1 I_g d\mu$, $F_g = \int_{-1}^1 \mu I_g d\mu$

$$\begin{aligned} \frac{\partial E_g}{\partial t} + \frac{\partial F_g}{\partial x} + c \kappa_g(T) E_g &= 4\pi \kappa_g(T) B_g(T), \\ \frac{1}{c} \frac{\partial F_g}{\partial t} + c \frac{\partial}{\partial x} (f_g[l] E_g) + \kappa_g(T) F_g &= 0 \end{aligned}$$

- Effective grey problem for $E = \sum_{g=1}^{N_g} E_g$, $F = \sum_{g=1}^{N_g} F_g$

$$\begin{aligned} \frac{\partial E}{\partial t} + \frac{\partial F}{\partial x} + c \bar{\kappa}_E E &= c \bar{\kappa}_{BaR} T^4 \\ \frac{1}{c} \frac{\partial F}{\partial t} + c \frac{\partial (\bar{f} E)}{\partial x} + \bar{\kappa}_R F + \bar{\eta} E &= 0 \\ \frac{\partial \varepsilon(T)}{\partial t} &= c (\bar{\kappa}_E E - \bar{\kappa}_{BaR} T^4) \end{aligned}$$

- POD Galerkin Projected Boltzmann transport equation

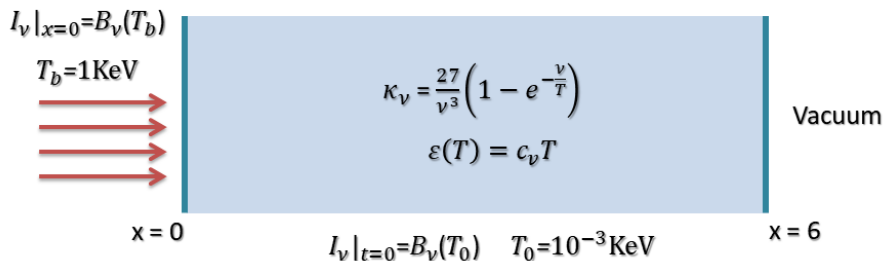
$$\frac{1}{c\Delta t^n} (\lambda_\ell^n - \lambda_\ell^{n-1}) + \sum_{\ell'=1}^r \lambda_{\ell'}^n \langle \mathbf{u}_\ell, \mathcal{L}_h \mathbf{u}_{\ell'} \rangle_W + \sum_{\ell'=1}^r \lambda_{\ell'}^n \langle \mathbf{u}_\ell, \mathcal{K}_h^n(T) \mathbf{u}_{\ell'} \rangle_W = \langle \mathbf{u}_\ell, \mathbf{Q}^n(T) \rangle_W$$

- Approximate intensities $\mathbf{I}_r^u = \sum_{\ell=1}^r \lambda_\ell^n \mathbf{u}_\ell \implies \tilde{\mathbf{f}}_g[\mathbf{I}_r^u]$
- Multigroup quasidiffusion equations

$$\begin{aligned} \frac{\partial E_g}{\partial t} + \frac{\partial F_g}{\partial x} + c\kappa_g(T)E_g &= 4\pi\kappa_g(T)B_g(T), \\ \frac{1}{c} \frac{\partial F_g}{\partial t} + c \frac{\partial}{\partial x} (\tilde{\mathbf{f}}_g[\mathbf{I}_r^u] E_g) + \kappa_g(T)F_g &= 0 \end{aligned}$$

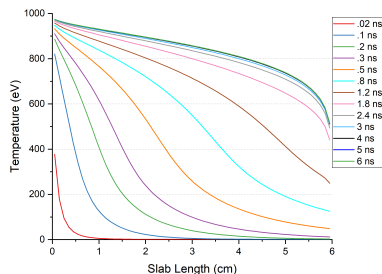
- Effective grey problem

$$\begin{aligned} \frac{\partial E}{\partial t} + \frac{\partial F}{\partial x} + c\bar{\kappa}_E E &= c\bar{\kappa}_{BaR} T^4 \\ \frac{1}{c} \frac{\partial F}{\partial t} + c \frac{\partial (\bar{\mathbf{f}} E)}{\partial x} + \bar{\kappa}_R F + \bar{\eta} E &= 0 \\ \frac{\partial \varepsilon(T)}{\partial t} &= c(\bar{\kappa}_E E - \bar{\kappa}_{BaR} T^4) \end{aligned}$$

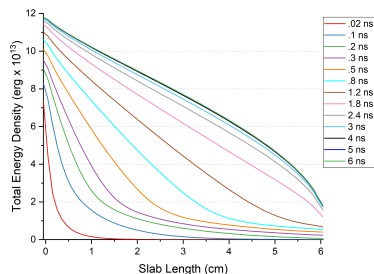


- Fleck & Cummings, 1971
- 17 frequency (energy) groups
- 60 spatial cells, $\Delta x = 0.1$ cm
- $\Delta t = 2 \times 10^{-2}$ ns
- $0 \leq t \leq 6$ ns, 300 time steps
- DS_4 Gaussian quadrature set
- Finite volume in space & fully implicit scheme for LOQD eqs.

- The F-C test is characterized by three distinct temporal stages
 - Rapid wave formation $t \in [0, 0.3\text{ns}]$
 - Propagation of wave from left to right $t \in (0.3, 1.2\text{ns}]$
 - Slow heating of entire domain towards steady state $t \in (1.2, 6\text{ns}]$



Temperature



Energy density

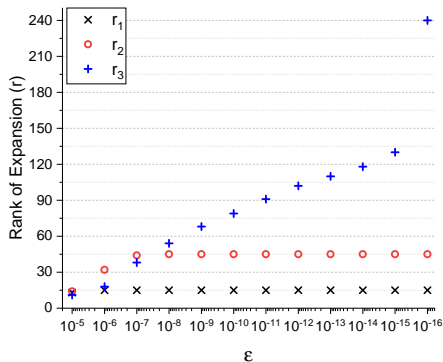
Calculation of Basis

- We calculate a unique POD basis for each distinct stage of the F-C test
- \mathbf{A}_1 , \mathbf{A}_2 , \mathbf{A}_3 with $d_i = \text{rank}(\mathbf{A}_i)$

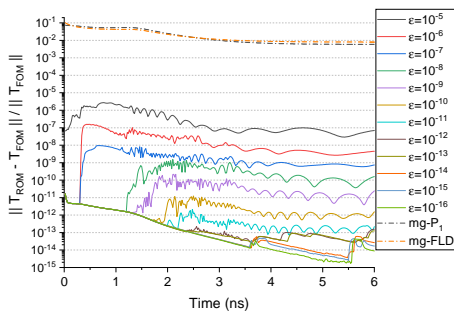
- ranks $r_i \leq d_i$ are calculated based off singular values of \mathbf{A}_i

$$\left(\frac{\sum_{\ell=r_i+1}^{d_i} \sigma_\ell^2}{\sum_{\ell=1}^{d_i} \sigma_\ell^2} \right)^{\frac{1}{2}} < \xi$$

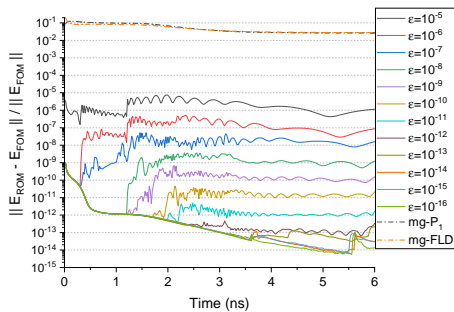
- Stage 1 (r_1): full rank = 15
- Stage 2 (r_2): full rank = 45
- Stage 3 (r_3): full rank = 240



- Relative errors in 2-norm at each instant of time

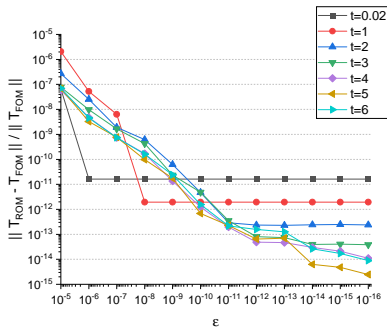


Temperature

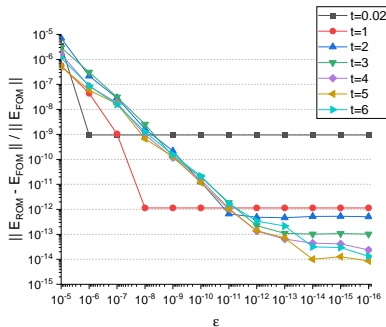


Energy density

- Relative errors in 2-norm plotted vs ξ for certain times



Temperature



Energy density

- We developed and tested a prototype advanced ROM for TRT in 1D geometry
- A POD Galerkin-Projected BTE is coupled with the low-order quasidiffusion equations to give approximate closure
- The Projected BTE is dense but contains many less degrees of freedom compared to the original BTE
- The expansion coefficients directly depend on material temperature, making the ROM naturally parametric
- The ROM was shown to produce highly accurate solutions, and converge uniformly to the reference solution as rank is increased
- Results are promising, and future work includes:
 - Extension of the method into 2D geometry
 - Investigation into performance over given parametric spaces